

# Chapter 1

## Fourier Series

Basic definitions and examples of Fourier series are given in Section 1. In Section 2 we prove the fundamental Riemann-Lebesgue lemma and discuss Fourier series from the mapping point of view. Pointwise and uniform convergence of the Fourier series of a function to the function itself under various regularity assumptions are studied in Section 3. As an application, it is shown that every continuous function can be approximated by polynomials in a uniform manner in Section 4. In the first appendix basic facts on series of functions are summarised. In the second appendix the concept of measure zero sets is discussed.

### 1.1 Definition and Examples

In Mathematical Analysis II power series have been studied. Now we come to Fourier series.

First of all, a **trigonometric series** is a series of functions of the form

$$\sum_{n=0}^{\infty} (a_n \cos nx + b_n \sin nx), \quad a_n, b_n \in \mathbb{R}.$$

As  $\cos 0x = 1$  and  $\sin 0x = 0$ , we always set  $b_0 = 0$  and express the series as

$$a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx).$$

It is called a **cosine series** if all  $b_n$ 's vanish and **sine series** if all  $a_n$ 's vanish. Trigonometric series form an important class of series of functions. In Mathematical Analysis II, we studied the convergence of the series of general functions. We recall

- Uniform convergence implies pointwise convergence of a series of functions,
- Absolute convergence implies pointwise convergence of a series of functions,
- Weierstrass M-Test for uniform and absolute convergence (see Appendix I).
- Uniform convergence preserves continuity.

For instance, using the fact  $|\cos nx|, |\sin nx| \leq 1$ , Weierstrass M-Test tells us that a trigonometric series is uniformly and absolutely convergent when its coefficients satisfy

$$\sum_n |a_n|, \quad \sum_n |b_n| < \infty ,$$

and this is the case when  $|a_n|, |b_n| \leq Cn^{-s}, \forall n \geq 1$ , for some constant  $C$  and  $s > 1$ . Since the partial sums are continuous functions and uniform convergence preserves continuity, the infinite series

$$\varphi(x) \equiv a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$$

is a continuous function on  $\mathbb{R}$ . Actually, it is also periodic of period  $2\pi$ . For, by pointwise convergence, we have

$$\begin{aligned} \varphi(x + 2\pi) &= \lim_{n \rightarrow \infty} \sum_{k=0}^n (a_k \cos(kx + 2k\pi) + b_k \sin(kx + 2k\pi)) \\ &= \lim_{n \rightarrow \infty} \sum_{k=0}^n (a_k \cos kx + b_k \sin kx) \\ &= \varphi(x), \end{aligned}$$

hence it is  $2\pi$ -periodic.

Recall that there is a special power series associated to a function which is smooth at a certain point. Indeed, it is given by the Taylor's series at this point. Let the point be  $x_0$  and  $f$  is smooth in an open interval containing  $x_0$ , this series is given by

$$\sum_{n=0}^{\infty} c_n (x - x_0)^n, \quad c_n = \frac{f^{(n)}(x_0)}{n!} .$$

Similarly, there is a trigonometric series associated to an integrable function. It is called the Fourier series of the function. Let us define it now.

Given a  $2\pi$ -periodic function which is Riemann integrable function  $f$  on  $[-\pi, \pi]$ , its

**Fourier series** or **Fourier expansion** is the trigonometric series given by

$$\begin{aligned} a_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(y) \cos ny \, dy, \quad n \geq 1 \\ b_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(y) \sin ny \, dy, \quad n \geq 1 \quad \text{and} \\ a_0 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} f(y) \, dy. \end{aligned} \tag{1.1}$$

Note that  $a_0$  is the average of the function over the interval. From this definition we gather two basic information. First, the Fourier series of a function involves the integration of the function over an interval, hence any modification of the values of the function over a subinterval, not matter how small it is, may change the Fourier coefficients  $a_n$  and  $b_n$ . This is unlike a power series which only depends on the local properties (derivatives of all order at a designated point). We may say a Fourier series depends on the global information but a power series only depends on local information. Second, recalling from the theory of Riemann integral, we know that two integrable functions which are equal almost everywhere have the same integral. We will see that the converse is also true, namely, two functions with the same Fourier series are equal almost everywhere. Therefore, the Fourier series of two such functions are the same. In particular, the Fourier series of a function is completely determined with its value on the open interval  $(-\pi, \pi)$ , regardless its values at the endpoints.

The motivation of the Fourier series comes from the belief that for a “nice function” of period  $2\pi$ , its Fourier series converges to the function itself. In other words, we have

$$f(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx), \quad \forall x \in \mathbb{R}, \tag{1.2}$$

for some  $a_n$ 's and  $b_n$ 's. At this point we do not know what these coefficients are. We claim that whenever this holds, they must be given by (1.1). A formal argument proceeds as follows. Multiply (1.2) by  $\cos mx$ ,  $m \geq 1$ , and then integrate over  $[-\pi, \pi]$ . Using the formulas

$$\begin{aligned} \int_{-\pi}^{\pi} \cos nx \cos mx \, dx &= \begin{cases} \pi, & n = m \\ 0, & n \neq m \end{cases}, \\ \int_{-\pi}^{\pi} \cos nx \sin mx \, dx &= 0, \text{ and} \\ \int_{-\pi}^{\pi} \cos nx \, dx &= \begin{cases} 2\pi, & n = 0 \\ 0, & n \neq 0 \end{cases}, \end{aligned}$$

we arrive at the expression of  $a_n$ ,  $n \geq 0$ , in (1.2). Similarly, by multiplying (1.2) by  $\sin mx$  and then integrate over  $[-\pi, \pi]$ , one obtain the expression of  $b_n$ , in (1.2) after using

$$\int_{-\pi}^{\pi} \sin nx \sin mx \, dx = \begin{cases} \pi, & n = m \\ 0, & n \neq m \end{cases}.$$

Of course, (1.2) arises from the hypothesis that every sufficiently nice function of period  $2\pi$  is equal to its Fourier expansion. The study of under which “nice conditions” this could happen is one of the main objects in the theory of Fourier series.

We can associate a Fourier series for any integrable function on  $[-\pi, \pi]$ . As the right hand side of (1.2) consists of  $2\pi$ -periodic functions, it is natural to extend its left hand side, that is, the function  $f$  itself, as a  $2\pi$ -periodic function. The extension is straightforward. First of all, the real line can be expressed as the disjoint union of intervals  $((2n-1)\pi, (2n+1)\pi]$ ,  $n \in \mathbb{Z}$ . Each number  $x$  belongs to one and exactly one such interval. Let  $\tilde{f}(x) = f(x - 2n\pi)$  where  $n$  is the unique integer satisfying  $(2n-1)\pi < x \leq (2n+1)\pi$ . It is clear that  $\tilde{f}$  is equal to  $f$  on  $(-\pi, \pi]$ . As the original function is defined on  $[-\pi, \pi]$ , apparently an extension in strict sense is possible only if  $f(-\pi) = f(\pi)$ . Since the function value at one point does not change the Fourier series, from now on it will be understood that the extension of a function to a  $2\pi$ -periodic function refers to the extension for the restriction of this function on  $(-\pi, \pi]$ . Note that for the  $2\pi$ -periodic extension of a continuous function on  $[-\pi, \pi]$  has a jump discontinuity at  $\pm\pi$  when  $f(\pi) \neq f(-\pi)$ . It is continuous on  $\mathbb{R}$  if and only if  $f(-\pi) = f(\pi)$ . In the following we will not distinguish  $f$  with its extension  $\tilde{f}$ .

We will use

$$f(x) \sim a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$$

to denote the fact that the right hand side of this expression is the Fourier series of  $f$ . Note that in general  $\sim$  cannot be replaced by  $=$ .

**Example 1.1** We consider the function  $f_1(x) = x$ . Its  $2\pi$ -periodic extension is a function smooth everywhere except jump discontinuities at  $(2n+1)\pi, n \in \mathbb{Z}$ . As  $f_1$  is odd and  $\cos nx$  is even,

$$\pi a_n = \int_{-\pi}^{\pi} x \cos nx \, dx = 0, \quad n \geq 0,$$

and

$$\begin{aligned} \pi b_n &= \int_{-\pi}^{\pi} x \sin nx \, dx \\ &= -x \frac{\cos nx}{n} \Big|_{-\pi}^{\pi} + \int_{-\pi}^{\pi} \frac{\cos nx}{n} \, dx \\ &= (-1)^{n+1} \frac{2\pi}{n}. \end{aligned}$$

Therefore,

$$f_1(x) \sim 2 \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} \sin nx.$$

Since  $f_1$  is an odd function, it is reasonable to see that no cosine functions are involved in its Fourier series. How about the convergence of this Fourier series? Although the

coefficients decay like  $O(1/n)$  as  $n \rightarrow \infty$ , its convergence is not clear at this moment. On the other hand, this Fourier series is equal to 0 at  $x = \pm\pi$  but  $f_1(\pm\pi) = \pi$ . So, one thing is sure, namely, the Fourier series is not always equal to its function. It is worthwhile to observe that the bad points  $\pm\pi$  are precisely the discontinuity points of  $f_1$ .

**Notation** The big  $O$  and small  $o$  notations are very convenient in analysis. We say a sequence  $\{x_n\}$  satisfies  $x_n = O(n^s)$  means that there exists a constant  $C$  independent of  $n$  such that  $|x_n| \leq Cn^s$ . When  $s$  is positive, it means the growth of  $\{x_n\}$  is not faster than  $n^s$ . When  $s$  is negative, the decay of  $\{x_n\}$  is not slower than  $n^s$ . On the other hand,  $x_n = o(n^s)$  means  $|x_n|/n^s \rightarrow 0$  as  $n \rightarrow \infty$ .

**Example 1.2** Next consider the function  $f_2(x) = x^2$ . Unlike the previous example, its  $2\pi$ -periodic extension is continuous on  $\mathbb{R}$ . (However it is no longer differentiable at  $(2n+1)\pi$ .) After performing integration by parts, the Fourier series of  $f_2$  is seen to be

$$f_2(x) \equiv x^2 \sim \frac{\pi^2}{3} - 4 \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n^2} \cos nx.$$

As  $f_2$  is an even function, this is a cosine series. The rate of decay of the Fourier coefficients is like  $O(1/n^2)$ . Using Weierstrass M-test, this series converges uniformly to a continuous function. Later we will see that this continuous function is equal to  $f_2$ , but at this stage we do not know.

We list more examples of Fourier series of functions and leave them for you to verify.

$$(a) f_3(x) \equiv |x| \sim \frac{\pi}{2} - \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2n-1)^2} \cos(2n-1)x,$$

$$(b) f_4(x) = \begin{cases} 1, & x \in [0, \pi] \\ -1, & x \in (-\pi, 0) \end{cases} \sim \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{2n-1} \sin(2n-1)x,$$

$$(c) f_5(x) = \begin{cases} x(\pi-x), & x \in [0, \pi] \\ x(\pi+x), & x \in (-\pi, 0) \end{cases} \sim \frac{8}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2n-1)^3} \sin(2n-1)x.$$

Let  $\{c_n\}_{-\infty}^{\infty}$  be a bisequence of complex numbers. (As contrast to a sequence of complex numbers which is a map from  $\mathbb{N}$  to  $\mathbb{C}$ , a bisequence is a map from  $\mathbb{Z}$  to  $\mathbb{C}$ .) A (complex) trigonometric series is the infinite series associated to the bisequence  $\{c_n e^{inx}\}_{-\infty}^{\infty}$  and is denoted by  $\sum_{-\infty}^{\infty} c_n e^{inx}$ . To be in line with the real case, it is said to be convergent at  $x$  if

$$\lim_{n \rightarrow \infty} \sum_{k=-n}^n c_k e^{ikx}$$

exists. Now, a complex Fourier series can be associated to a complex-valued function. Let  $f$  be a  $2\pi$ -periodic complex-valued function which is integrable on  $[-\pi, \pi]$ . Its Fourier series is given by the series

$$f(x) \sim \sum_{n=-\infty}^{\infty} c_n e^{inx},$$

where the Fourier coefficients  $c_n$  are defined to be

$$c_n = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) e^{-inx} dx, \quad n \in \mathbb{Z}.$$

Here for a complex-valued function  $f$ , its integration over some  $[a, b]$  is defined to be

$$\int_a^b f(x) dx \equiv \int_a^b f_1(x) dx + i \int_a^b f_2(x) dx,$$

where  $f_1$  and  $f_2$  are respectively the real and imaginary parts of  $f$ . And differentiation is understood as

$$f' \equiv f'_1 + i f'_2.$$

It is called integrable/differentiable if both real and imaginary parts are integrable/differentiable. The same as in the real case, formally the expression of  $c_n$  is obtained as in the real case by first multiplying the relation

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx}$$

with  $e^{imx}$  and then integrating over  $[-\pi, \pi]$  with the help from the relation

$$\int_{-\pi}^{\pi} e^{imx} e^{-inx} dx = \begin{cases} 2\pi, & n = m \\ 0, & n \neq m \end{cases}.$$

When  $f$  is of real-valued, there are two Fourier series associated to  $f$ , that is, the real and the complex ones. To relate them it is enough to observe the Euler's formula  $e^{i\theta} = \cos \theta + i \sin \theta$ , so for  $n \geq 1$

$$\begin{aligned} 2\pi c_n &= \int_{-\pi}^{\pi} f(x) e^{-inx} dx \\ &= \int_{-\pi}^{\pi} f(x) (\cos nx - i \sin nx) dx \\ &= \int_{-\pi}^{\pi} f(x) \cos nx dx - i \int_{-\pi}^{\pi} f(x) \sin nx dx \\ &= \pi(a_n - ib_n). \end{aligned}$$

we see that

$$c_n = \frac{1}{2}(a_n - ib_n), \quad n \geq 1, \quad c_0 = a_0.$$

By a similar computation, we have

$$c_n = \frac{1}{2}(a_{-n} + ib_{-n}), \quad n \leq -1.$$

It follows that  $c_{-n} = \overline{c_n}$  for all  $n$ . The complex form of Fourier series sometimes makes expressions and computations more elegant. We will use it whenever it makes things simpler.

We have been working on the Fourier series of  $2\pi$ -periodic functions. For functions of  $T$ -period, their Fourier series are not the same. They can be found by a scaling argument. Let  $f$  be  $T$ -periodic. The function  $g(x) = f(Tx/2\pi)$  is a  $2\pi$ -periodic function. Thus,

$$f\left(\frac{Tx}{2\pi}\right) = g(x) \sim a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx),$$

where  $a_0, a_n, b_n, n \geq 1$  are the Fourier coefficients of  $g$ . By a change of variables, we can express everything inside the coefficients in terms of  $f$ ,  $\cos 2n\pi x/T$  and  $\sin 2n\pi x/T$ . The result is

$$f(x) \sim a_0 + \sum_{n=1}^{\infty} \left( a_n \cos \frac{2n\pi}{T}x + b_n \sin \frac{2n\pi}{T}x \right),$$

where

$$\begin{aligned} a_n &= \frac{2}{T} \int_{-T/2}^{T/2} f(y) \cos \frac{2n\pi}{T}y \, dy, \\ b_n &= \frac{2}{T} \int_{-T/2}^{T/2} f(y) \sin \frac{2n\pi}{T}y \, dy, \quad n \geq 1, \quad \text{and} \\ a_0 &= \frac{1}{T} \int_{-T/2}^{T/2} f(y) \, dy. \end{aligned}$$

Here the integrals could be replaced by any interval of length  $T$ . It reduces to (1.1) when  $T$  is equal to  $2\pi$ . From now on, when we talk about the Fourier series of a function on an interval of length  $T$ , it is understood that the “building blocks” are  $\cos 2n\pi/Tx$  and  $\sin 2n\pi/Tx, n \geq 1$  or  $e^{2n\pi ix}, n \in \mathbb{Z}$ .

## 1.2 Riemann-Lebesgue Lemma

From the examples of Fourier series of functions in the previous section we see that the coefficients decay to 0 eventually. We will show that this is generally true. This is the content of the following result.

**Theorem 1.1 (Riemann-Lebesgue Lemma).** For  $f \in R[a, b]$ ,

$$\int_a^b f(x) \cos nx \, dx, \quad \int_a^b f(x) \sin nx \, dx \rightarrow 0, \quad \text{as } n \rightarrow \infty.$$

In particular, taking  $[a, b] = [-\pi, \pi]$ , the Fourier coefficients of  $f$ ,  $a_n, b_n \rightarrow 0$  as  $n \rightarrow \infty$ .

To prepare for the proof, we examine how to approximate an integrable function by step functions. Let  $a_0 = a < a_1 < \dots < a_N = b$  be a partition of  $[a, b]$ . A **step function**  $s$  satisfies  $s(x) = s_j, \forall x \in (a_j, a_{j+1}], \forall j \geq 0$ . The value of  $s$  at  $a$  is not important, but for definiteness let's set  $s(a) = s_0$ . We can express a step function in a better form by introducing the **characteristic function**  $\chi_E$  of a set  $E \subset \mathbb{R}$ :

$$\chi_E(x) = \begin{cases} 1, & x \in E, \\ 0, & x \notin E. \end{cases}$$

Then,

$$s(x) = \sum_{j=0}^{N-1} s_j \chi_{I_j}, \quad I_j = (a_j, a_{j+1}], \quad j \geq 1, \quad I_0 = [a_0, a_1].$$

**Lemma 1.2.** For every step function  $s$ , there exists some constant  $C$  independent of  $n$  such that

$$\left| \int_a^b s(x) \cos nx \, dx \right|, \quad \left| \int_a^b s(x) \sin nx \, dx \right| \leq \frac{C}{n}, \quad \forall n \geq 1.$$

*Proof.* Let  $s(x) = \sum_{j=0}^{N-1} s_j \chi_{I_j}$ . We have

$$\begin{aligned} \int_a^b s(x) \cos nx \, dx &= \int_a^b \sum_{j=0}^{N-1} s_j \chi_{I_j} \cos nx \, dx \\ &= \sum_{j=0}^{N-1} s_j \int_{a_j}^{a_{j+1}} \cos nx \, dx \\ &= \frac{1}{n} \sum_{j=0}^{N-1} s_j (\sin na_{j+1} - \sin na_j). \end{aligned}$$

It follows that

$$\left| \int_a^b s(x) \cos nx \, dx \right| \leq \frac{C}{n}, \quad \forall n \geq 1, \quad C = 2 \sum_{j=0}^{N-1} |s_j|.$$

Clearly a similar estimate holds for the other case. □

**Lemma 1.3.** *Let  $f \in R[a, b]$ . Given  $\varepsilon > 0$ , there exists a step function  $s$  such that  $s \leq f$  on  $[a, b]$  and*

$$0 \leq \int_a^b (f - s) < \varepsilon.$$

*Proof.* As  $f$  is integrable, it can be approximated from below by its Darboux lower sums. In other words, for  $\varepsilon > 0$ , we can find a partition  $a = a_0 < a_1 < \cdots < a_N = b$  such that

$$0 \leq \int_a^b f - \sum_{j=0}^{N-1} m_j (a_{j+1} - a_j) < \varepsilon,$$

where  $m_j = \inf \{f(x) : x \in [a_j, a_{j+1}]\}$ . It follows that

$$0 \leq \int_a^b (f - s) < \varepsilon$$

after setting

$$s(x) = \sum_{j=0}^{N-1} m_j \chi_{I_j}, \quad I_j = (a_j, a_{j+1}], \quad j \geq 1, \quad I_0 = [a_0, a_1].$$

□

Now we prove Theorem 1.1. For  $\varepsilon > 0$ , we can find  $s$  as constructed in Lemma 1.3 such that  $0 \leq f - s$  and

$$0 \leq \int_a^b (f - s) < \frac{\varepsilon}{2}.$$

By Lemma 1.2, there exists some  $n_0$  such that

$$\left| \int_a^b s(x) \cos nx dx \right| < \frac{\varepsilon}{2},$$

for all  $n \geq n_0$ . Therefore,

$$\begin{aligned} \left| \int_a^b f(x) \cos nx dx \right| &\leq \left| \int_a^b (f - s) \cos nx dx \right| + \left| \int_a^b s(x) \cos nx dx \right| \\ &\leq \int_a^b |f - s| + \frac{\varepsilon}{2} \\ &< \frac{\varepsilon}{2} + \frac{\varepsilon}{2} = \varepsilon. \end{aligned}$$

The same argument applies when  $\cos nx$  is replaced by  $\sin nx$ . The proof of Riemann-Lebesgue Lemma is completed.

There are different and simpler proofs of this lemma when  $f$  is differentiable or continuous, see exercise.

It is useful to bring in a “mapping” point of view between functions and their Fourier series. Let  $R_{2\pi}$  be the collection of all  $2\pi$ -periodic complex-valued functions (or real-valued functions depending the context) integrable on  $[-\pi, \pi]$  and  $\mathcal{C}$  consisting of all complex-valued bisequences  $\{c_n\}$  satisfying  $c_n \rightarrow 0$  as  $n \rightarrow \pm\infty$ . The Fourier series sets up a mapping  $\Phi$  from  $R_{2\pi}$  to  $\mathcal{C}$  by sending  $f$  to  $\{\hat{f}(n)\}$  where, to make things clear, we have let  $\hat{f}(n) = c_n$ , the  $n$ -th Fourier coefficient of  $f$ . When real-valued functions are considered, restricting to the subspace of  $\mathcal{C}$  given by those satisfying  $c_{-n} = \overline{c_n}$ ,  $\Phi$  maps all real-valued functions into this subspace. Alternatively, one may consider the mapping from the space of integrable functions to the space consisting of two sequences of real numbers  $\{a_0, a_1, a_2, \dots; b_1, b_2, \dots\}$ . In the following discussion we shall focus on the complex case, and let you fill in the real case.

With this correspondence at hand, our first question is: Is  $\Phi$  one-to-one? Clearly the answer is no, for two functions which differ on a set of measure zero have the same Fourier coefficients. However, we will establish the following result:

**Uniqueness Theorem.** *The Fourier series of two integrable functions in  $R_{2\pi}$  coincide if and only if they are equal almost everywhere. The Fourier series of two continuous  $2\pi$ -periodic functions coincide if and only if they are identical.*

Thus  $\Phi$  is essentially one-to-one. We may also study how various properties in  $R_{2\pi}$  and  $\mathcal{C}$  correspond under  $\Phi$ . In fact, there are obvious and surprising ones. Some of them are listed below and more can be found in the exercise. Observe that both  $R_{2\pi}$  and  $\mathcal{C}$  carry the structure of a vector space over  $\mathbb{C}$ .

**Property 1.**  $\Phi$  is a linear map. Observe that both  $R_{2\pi}$  and  $\mathcal{C}$  form vector spaces over  $\mathbb{R}$  or  $\mathbb{C}$ . The linearity of  $\Phi$  is clear from its definition.

**Property 2.** When  $f \in R_{2\pi}$  is  $k$ -th differentiable on  $\mathbb{R}$  and all derivatives up to  $k$ -th order belong to  $R_{2\pi}$ ,  $\hat{f}^k(n) = (in)^k \hat{f}(n)$  for all  $n \in \mathbb{Z}$ . See Proposition 1.4 below for a proof. This property shows that differentiation turns into the multiplication of a factor  $(in)^k$  under  $\Phi$ . This is amazing!

**Property 3.** Every translation in  $\mathbb{R}$  induces a “translation operation” on functions defined on  $\mathbb{R}$ . More specifically, for  $a \in \mathbb{R}$ , set  $f_a(x) = f(x - a)$ ,  $x \in \mathbb{R}$ . Clearly  $f_a$  belongs to  $R_{2\pi}$ . We have  $\hat{f}_a(n) = e^{-ina} \hat{f}(n)$ . This property follows directly from the definition. It shows that a translation in  $R_{2\pi}$  turns into the multiplication of a factor  $e^{-ina}$  under  $\Phi$ , or, by a rotation of angle  $-na/2\pi$ .

**Proposition 1.4.** *Let  $f$  be a continuous  $2\pi$ -periodic function which is differentiable on  $[-\pi, \pi]$  with  $f' \in R_{2\pi}$ . Then*

$$\hat{f}'(n) = in\hat{f}(n).$$

When  $f$  is real-valued,

$$f'(x) \sim \sum_{n=1}^{\infty} (nb_n \cos nx - na_n \sin nx),$$

where

$$f(x) \sim a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx) .$$

*Proof.* Let  $f' \sim \sum_{-\infty}^{\infty} c'_n e^{inx}$ . We have

$$\begin{aligned} c'_n &\equiv \int f'(x) e^{-inx} dx \\ &= \int (f'_1(x) + if'_2(x))(\cos nx - i \sin nx) dx \\ &= \int (f'_1(x) \cos nx + f'_2(x) \sin nx) dx + i \int (-f'_1(x) \sin nx + f'_2(x) \cos nx) dx \\ &= n \int (f_1(x) \sin nx - f_2(x) \cos nx) dx + in \int (f_1(x) \cos nx + f_2(x) \sin nx) dx \\ &= n \int (f_1(x) + if_2(x)) \sin nxdx + in \int (f_1(x) + if_2(x)) \cos nxdx \\ &= ni \int (f_1(x) + if_2(x))(\cos nx - i \sin nx) dx \\ &\equiv nic_n . \end{aligned}$$

Here the integration is from  $-\pi$  to  $\pi$ . Alternatively, one may verify that the integration by parts formula  $fg' = f'g + fg'$  also holds for complex-valued  $f$  and  $g$ , thus

$$\begin{aligned} c'_n &\equiv \int_{-\pi}^{\pi} f'(x) e^{-inx} dx \\ &= f(x) e^{-inx} \Big|_{-\pi}^{\pi} - \int_{-\pi}^{\pi} f(x) (-inx) e^{-inx} dx \\ &= inc_n . \end{aligned}$$

In the real case, let the Fourier coefficients of  $f'$  be  $a'_n$  and  $b'_n$ 's.

$$\begin{aligned} \pi a'_n &= \int_{-\pi}^{\pi} f'(y) \cos ny dy \\ &= f(y) \cos ny \Big|_{-\pi}^{\pi} - \int_{-\pi}^{\pi} f(y) (-n \sin ny) dy \\ &= n \int_{-\pi}^{\pi} f(y) \sin ny dy \\ &= \pi n b_n . \end{aligned}$$

Similarly,

$$\begin{aligned}
 \pi b'_n &= \int_{-\pi}^{\pi} f'(y) \sin ny \, dy \\
 &= f(y) \sin ny \Big|_{-\pi}^{\pi} - \int_{-\pi}^{\pi} f(y) n \cos ny \, dy \\
 &= -n \int_{-\pi}^{\pi} f(y) \cos ny \, dy \\
 &= -\pi n a_n.
 \end{aligned}$$

□

Property 2 links the regularity of the function to the rate of decay of its Fourier coefficients. This is an extremely important property. When  $f$  is a  $2\pi$ -periodic function whose derivatives up to  $k$ -th order belong to  $R_{2\pi}$ , applying Riemann-Lebesgue lemma to  $f^{(k)}$  we know that  $\hat{f}^{(k)}(n) = o(1)$  as  $n \rightarrow \infty$ . By Property 2 it follows that  $\hat{f}(n) = o(n^{-k})$ , that is, the Fourier coefficients of  $f$  decay faster than  $n^{-k}$ . Since  $\sum_{n=1}^{\infty} n^{-2} < \infty$ , an application of Weierstrass M-test establishes the following result: The Fourier series of  $f$  converges uniformly provided  $f, f'$  and  $f''$  belong to  $R_{2\pi}$ . Therefore, the function

$$g(x) \equiv a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx) ,$$

is a continuous  $2\pi$ -periodic function. Using its uniform convergence, we see that the Fourier coefficients of  $g$  are given by  $a_n$  and  $b_n$ , the same as  $f$ . By the Uniqueness Theorem, we conclude that  $g$  is equal to  $f$  almost everywhere. And since these two functions are continuous, they must be equal everywhere (an easy exercise). We conclude that the Fourier series of  $f$  is equal to  $f$  provided  $f, f', f'' \in R_{2\pi}$ . A more general result will be proved in the next section.

### 1.3 Convergence of Fourier Series

In this section we study the convergence of the Fourier series of a function to the function itself. Recall that the series  $a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$ , or  $\sum_{n=-\infty}^{\infty} c_n e^{inx}$ , where  $a_n, b_n, c_n$  are the Fourier coefficients of a function  $f$  converges to  $f$  at  $x$  means that the  $n$ -th partial sum

$$(S_n f)(x) = a_0 + \sum_{k=1}^n (a_k \cos kx + b_k \sin kx) , \quad S_0 f = a_0 ,$$

or

$$(S_n f)(x) = \sum_{k=-n}^n c_k e^{ikx}$$

converges to  $f(x)$  as  $n \rightarrow \infty$ .

We start by expressing the partial sums in closed form. Indeed,

$$\begin{aligned}
 (S_n f)(x) &= a_0 + \sum_{k=1}^n (a_k \cos kx + b_k \sin kx) \\
 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} f + \sum_{k=1}^n \frac{1}{\pi} \int_{-\pi}^{\pi} f(y) (\cos ky \cos kx + \sin ky \sin ky) dy \\
 &= \frac{1}{\pi} \int_{-\pi}^{\pi} \left( \frac{1}{2} + \sum_{k=1}^n \cos k(y-x) \right) f(y) dy \\
 &= \frac{1}{\pi} \int_{x-\pi}^{x+\pi} \left( \frac{1}{2} + \sum_{k=1}^n \cos kz \right) f(x+z) dz \\
 &= \frac{1}{\pi} \int_{-\pi}^{\pi} \left( \frac{1}{2} + \sum_{k=1}^n \cos kz \right) f(x+z) dz,
 \end{aligned}$$

where in the last step we have used the fact that the integrals over any two periods are the same. Using the elementary formula

$$\cos \theta + \cos 2\theta + \cdots + \cos n\theta = \frac{\sin \left( n + \frac{1}{2} \right) \theta - \sin \frac{1}{2} \theta}{2 \sin \frac{\theta}{2}}, \quad \theta \neq 0,$$

we obtain

$$\frac{1}{2} + \sum_{k=1}^n \cos k\theta = \frac{\sin \left( n + \frac{1}{2} \right) \theta}{2 \sin \frac{\theta}{2}}.$$

Noting that by the L'Hospital Rule,

$$\lim_{\theta \rightarrow 0} \frac{\sin \left( n + \frac{1}{2} \right) \theta}{2 \sin \frac{\theta}{2}} = \frac{2n+1}{2},$$

we introduce the **Dirichlet kernel**  $D_n$  by

$$D_n(x) = \begin{cases} \frac{\sin \left( n + \frac{1}{2} \right) x}{2\pi \sin \frac{1}{2} x}, & x \neq 0 \\ \frac{2n+1}{2\pi}, & x = 0. \end{cases}$$

(In fact, there are infinitely many Dirichlet kernels indexed by  $n$ , but usually people refer them as one.) It is a continuous,  $2\pi$ -periodic function. We have successfully expressed the partial sums of the Fourier series in the following closed form:

$$(S_n f)(x) = \int_{-\pi}^{\pi} D_n(z) f(x+z) dz,$$

Taking  $f \equiv 1$ , we have  $S_n f = 1$  for all  $n$ . Hence

$$1 = \int_{-\pi}^{\pi} D_n(z) dz.$$

We have arrived at the fundamental relation

$$(S_n f)(x) - f(x) = \int_{-\pi}^{\pi} D_n(z)(f(x+z) - f(x)) dz. \quad (1.3)$$

In order to show  $S_n f(x) \rightarrow f(x)$ , it suffices to show the right hand side of (1.3) tends to 0 as  $n \rightarrow \infty$ .

The Dirichlet kernel plays a crucial role in the study of the convergence of Fourier series. We list some of its properties as follows.

**Property I.**  $D_n(z)$  is an even, continuous,  $2\pi$ -periodic function vanishing at  $z = 2k\pi/(2n+1)$ ,  $-n \leq k \leq n$ , on  $[-\pi, \pi]$ .

**Property II.**  $D_n$  attains its maximum value  $(2n+1)/2\pi$  at 0.

**Property III.**

$$\int_{-\pi}^{\pi} D_n(z) dz = 1.$$

**Property IV.** For every  $\delta > 0$ ,

$$\int_0^{\delta} |D_n(z)| dz \rightarrow \infty, \quad \text{as } n \rightarrow \infty.$$

Only the last property needs a proof. Indeed, for each  $n$  we can fix an  $N$  such that  $\pi N < (2n+1)\delta/2 \leq (N+1)\pi$ , so  $N \rightarrow \infty$  as  $n \rightarrow \infty$ . We compute

$$\begin{aligned} \int_0^{\delta} |D_n(z)| dz &= \int_0^{\delta} \frac{|\sin(n + \frac{1}{2})z|}{2\pi |\sin \frac{z}{2}|} dz \\ &\geq \int_0^{(n+\frac{1}{2})\delta} \frac{|\sin t|}{2\pi t/2} dt \quad (\text{use } t/2 \geq \sin(t/2)) \\ &\geq \frac{1}{\pi} \int_0^{N\pi} \frac{|\sin t|}{t} dt \\ &= \frac{1}{\pi} \sum_{k=1}^N \int_{(k-1)\pi}^{k\pi} \frac{|\sin t|}{t} dt \\ &\geq \frac{1}{\pi} \sum_{k=1}^N \int_0^{\pi} \frac{|\sin s|}{k\pi} ds \quad (\text{use } 1/t \geq 1/(k\pi) \text{ on } [(k-1)\pi, k\pi]) \\ &= \frac{2}{\pi^2} \sum_{k=1}^N \frac{1}{k}, \quad \text{as } \int_0^{\pi} |\sin s| ds = 2, \\ &\rightarrow \infty, \end{aligned}$$

as  $N \rightarrow \infty$ .

To elucidate the effect of the kernel, we fix a small  $\delta > 0$  and introduce the bump function  $\Phi_\delta$  which satisfies (a)  $\Phi_\delta \in C(\mathbb{R})$ ,  $\Phi_\delta \equiv 0$  outside  $(-\delta, \delta)$ , (b)  $0 < \Phi_\delta \leq 1$  on  $(-\delta, \delta)$  and (c)  $\Phi_\delta = 1$  on  $(-\delta/2, \delta/2)$ . We split the integral into two parts:

$$\int_{-\pi}^{\pi} \Phi_\delta(z) D_n(z) (f(x+z) - f(x)) dz,$$

and

$$\int_{-\pi}^{\pi} (1 - \Phi_\delta(z)) D_n(z) (f(x+z) - f(x)) dz .$$

The second integral can be written as

$$\int_E \frac{(1 - \Phi_\delta(z))(f(x+z) - f(x))}{2\pi \sin \frac{z}{2}} (\sin \frac{z}{2} \cos nz + \cos \frac{z}{2} \sin nz) dz ,$$

where  $E$  is the union of two intervals  $[-\pi, -\delta/2]$  and  $[\delta/2, \pi]$ . As  $|\sin z/2|$  has a positive lower bound on  $E$  (in fact,  $|\sin z/2| \geq \sin \delta/2 > 0$  for  $\delta \in (0, \pi/2)$ ), both the functions

$$\frac{1 - \Phi_\delta(z)(f(x+z) - f(x)) \sin z/2}{2\pi \sin \frac{z}{2}}$$

and

$$\frac{1 - \Phi_\delta(z)(f(x+z) - f(x)) \cos z/2}{2\pi \sin \frac{z}{2}}$$

and belong to  $R[E]$  and the second integral tends to 0 as  $n \rightarrow \infty$  in view of Riemann-Lebesgue lemma. The trouble lies on the first integral. It can be estimated by

$$\int_{-\delta}^{\delta} |D_n(z)| |f(x+z) - f(x)| dz.$$

Unfortunately, in view of Property IV, no matter how small  $\delta$  is, this term may go to  $\infty$  so it is not clear how to estimate this integral.

The difficulty can be resolved by imposing a further regularity assumption on the function. First a definition. A function  $f$  defined on  $[a, b]$  is called **Lipschitz continuous** at  $x \in [a, b]$  if there exist  $L$  such that

$$|f(y) - f(x)| \leq L |y - x|, \quad \forall y \in [a, b]. \quad (1.4)$$

**Theorem 1.5.** *Let  $f$  be a  $2\pi$ -periodic function integrable on  $[-\pi, \pi]$ . Suppose that  $f$  is Lipschitz continuous at  $x$ . Then  $\{S_n f(x)\}$  converges to  $f(x)$  as  $n \rightarrow \infty$ .*

*Proof.* Let  $\Phi_\delta$  be the bump function as before. We write

$$\begin{aligned}
(S_n f)(x) - f(x) &= \int_{-\pi}^{\pi} D_n(z)(f(x+z) - f(x)) dz \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} \frac{\sin(n + \frac{1}{2})z}{\sin \frac{z}{2}} (f(x+z) - f(x)) dz \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} \Phi_\delta(z) \frac{\sin(n + \frac{1}{2})z}{\sin \frac{z}{2}} (f(x+z) - f(x)) dz \\
&\quad + \frac{1}{\pi} \int_{-\pi}^{\pi} (1 - \Phi_\delta(z)) \frac{\sin(n + \frac{1}{2})z}{\sin \frac{z}{2}} (f(x+z) - f(x)) dz \\
&\equiv I + II .
\end{aligned}$$

By our assumption on  $f$ ,

$$|f(x+z) - f(x)| \leq L|z| .$$

Using  $\sin \theta / \theta \rightarrow 1$  as  $\theta \rightarrow 0$ , there exists  $\delta$  such that  $2|\sin z/2| \geq |z/2|$  for all  $z, |z| < \delta$ . For  $z, |z| < \delta$ , we have  $|f(x+z) - f(x)|/|\sin z/2| \leq 4L$  and

$$\begin{aligned}
|I| &\leq \frac{1}{2\pi} \int_{-\delta}^{\delta} \Phi_\delta(z) \frac{|\sin(n + \frac{1}{2})z|}{|\sin \frac{z}{2}|} |f(x+z) - f(x)| dz \\
&\leq \frac{1}{2\pi} \int_{-\delta}^{\delta} 4L dz \\
&= \frac{4\delta L}{\pi} .
\end{aligned} \tag{1.5}$$

For  $\varepsilon > 0$ , we further restrict and fix a  $\delta$  so that

$$\frac{4\delta L}{\pi} < \frac{\varepsilon}{2} . \tag{1.6}$$

After fixing  $\delta$ , we turn to the second integral

$$\begin{aligned}
II &= \frac{1}{2\pi} \int_{-\pi}^{\pi} \frac{(1 - \Phi_\delta(z))(f(x+z) - f(x))}{\sin \frac{z}{2}} \sin(n + \frac{1}{2})z dz \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} \frac{(1 - \Phi_\delta(z))(f(x+z) - f(x))}{\sin \frac{z}{2}} \left( \sin \frac{z}{2} \cos nz + \cos \frac{z}{2} \sin nz \right) dz \\
&\equiv \int_E F_1(z, x) \sin nz dz + \int_E F_2(z, x) \cos nz dz,
\end{aligned}$$

where  $E = [-\pi, -\delta/2] \cup [\delta/2, \pi]$ ,

$$F_1(z, x) = \frac{1}{2\pi} \frac{(1 - \Phi_\delta(z))(f(x+z) - f(x))}{\sin \frac{z}{2}} \sin \frac{z}{2} ,$$

and

$$F_2(z, x) = \frac{1}{2\pi} \frac{(1 - \Phi_\delta(z))(f(x+z) - f(x))}{\sin \frac{z}{2}} \cos \frac{z}{2} .$$

As  $1 - \Phi_\delta(z) = 0$  on  $[-\delta/2, \delta/2]$ , these two functions vanish outside the two intervals  $[-\pi, -\delta/2]$  and  $[\delta/2, \pi]$ . Now  $|\sin(z/2)|$  has a positive lower bound on  $E$ , so  $F_1$  and  $F_2$  are integrable on  $[-\pi, \pi]$ . By Riemann-Lebesgue Lemma, for  $\varepsilon > 0$ , there is some  $n_0$  such that

$$\left| \int_{-\pi}^{\pi} F_1(z, x) \cos nz \, dz \right|, \left| \int_{-\pi}^{\pi} F_2(z, x) \sin nz \, dz \right| < \frac{\varepsilon}{4}, \quad \forall n \geq n_0. \quad (1.7)$$

Putting (1.5), (1.6) and (1.7) together,

$$|S_n f(x) - f(x)| < \frac{\varepsilon}{2} + \frac{\varepsilon}{4} + \frac{\varepsilon}{4} = \varepsilon, \quad \forall n \geq n_0.$$

We have shown that  $S_n f(x)$  tends to  $f(x)$  whenever  $f$  is Lipschitz continuous at  $x$ .  $\square$

A careful examination of it reveals a convergence result for functions with jump discontinuity after using the evenness of the Dirichlet kernel.

**Theorem 1.6.** *Let  $f$  be a  $2\pi$ -periodic function integrable on  $[-\pi, \pi]$ . Suppose at some  $x \in [-\pi, \pi]$ ,  $\lim_{y \rightarrow x^+} f(y)$  and  $\lim_{y \rightarrow x^-} f(y)$  exist and there are  $\delta > 0$  and constant  $L$  such that*

$$|f(y) - f(x^+)| \leq L(y - x), \quad \forall y, \quad x < y < \pi,$$

and

$$|f(y) - f(x^-)| \leq L(x - y), \quad \forall y, \quad -\pi < y < x.$$

Then  $\{S_n f(x)\}$  converges to  $(f(x^+) + f(x^-))/2$  as  $n \rightarrow \infty$ .

Here  $f(x^+)$  and  $f(x^-)$  stand for  $\lim_{y \rightarrow x^+} f(y)$  and  $\lim_{y \rightarrow x^-} f(y)$  respectively.

The proof of this theorem starts with the observation

$$\int_0^\pi D_n(z) \, dz = \frac{1}{2}, \quad \int_{-\pi}^0 D_n(z) \, dz = \frac{1}{2},$$

which is due to the fact that  $D_n$  is an even function and

$$\int_{-\pi}^{\pi} D_n(z) \, dz = 1.$$

We write

$$S_n(x) - \frac{1}{2}(f(x^+) + f(x^-)) = \int_0^\pi D_n(z)(f(x+z) - f(x^+)) \, dz + \int_{-\pi}^0 D_n(z)(f(x+z) - f(x^-)) \, dz,$$

and apply the arguments in the proof of Theorem 1.5 to these two integrals separately.

A function  $f$  defined on  $[a, b]$  is called **uniformly Lipschitz continuous** if there exists an  $L$  such that

$$|f(x) - f(y)| \leq L|x - y|, \quad \forall x, y \in [a, b].$$

In other words,  $L$  is independent of  $x$ . Every uniformly Lipschitz continuous function is Lipschitz continuous at every point. Every continuously differentiable function on  $[a, b]$  is uniformly Lipschitz continuous. In fact, by the fundamental theorem of calculus, for  $x, y \in [a, b]$ ,

$$\begin{aligned} |f(y) - f(x)| &= \left| \int_x^y f'(t) dt \right| \\ &\leq M|y - x|, \end{aligned}$$

where  $M = \sup\{|f'(t)| : t \in [a, b]\}$ .

Now, we have a theorem on the uniform convergence of the Fourier series of a function to the function itself.

**Theorem 1.7.** *Let  $f$  be a  $2\pi$ -periodic function satisfying the Lipschitz condition on  $\mathbb{R}$ . Its Fourier series converges to itself uniformly as  $n \rightarrow \infty$ .*

In particular, it means that the Fourier series of a continuously differentiable  $2\pi$ -periodic function converges uniformly to itself.

*Proof.* Observe that when  $f$  is Lipschitz continuous on  $\mathbb{R}$ ,  $L$  is independent of  $x$  and (1.5), (1.6) hold uniformly in  $x$ . Thus the theorem follows if  $n_0$  in (1.7) can be chosen independent of  $x$ . This is the content of the lemma below. We apply it by taking  $F(s, t)$  to be  $F_1(z, x)$  or  $F_2(z, x)$  with  $s, t$  replaced by  $z, x$  respectively. □

**Lemma 1.8.** *Let  $F(s, t)$  be continuous in  $[a, b] \times [c, d]$ . Then*

$$g(n, t) \equiv \int_a^b F(s, t) e^{-ins} ds \rightarrow 0$$

*uniformly on  $[c, d]$  as  $n \rightarrow \infty$ .*

*Proof.* We need to show that for every  $\varepsilon > 0$ , there exists some  $n_0$  independent of  $t$  such that

$$|g(n, t)| < \varepsilon, \quad \forall n \geq n_0.$$

First of all, by a theorem in 2050, a continuous function on  $[a, b] \times [c, d]$  is uniformly continuous. For  $\varepsilon > 0$ , there is some  $\delta$  such that  $|F(s, t) - F(s', t')| < \varepsilon$  whenever  $|s - s'|, |t - t'| < \delta$ . Now, given  $\varepsilon > 0$ , we divide  $[c, d]$  into  $M$  many subintervals  $[t_j, t_{j+1}]$

of length less than  $\delta$ . For any  $t \in [c, d]$ , there is some  $t_j$  such that  $|t - t_j| < \delta$ , thus  $|F(s, t) - F(s, t_j)| < \varepsilon$ . Applying Riemann-Lebesgue lemma to the function  $F(s, t_j)$ , we find  $n_j$  such that

$$\left| \int_a^b F(s, t_j) e^{-ins} ds \right| < \varepsilon, \quad \forall n \geq n_j.$$

Therefore, for  $t \in [c, d]$ , and  $n \geq n_1, n_2, \dots, n_M$ ,

$$\begin{aligned} \left| \int_a^b F(s, t) e^{-ins} ds \right| &\leq \left| \int_a^b F(s, t_j) e^{-ins} ds \right| + \left| \int_a^b (F(s, t) - F(s, t_j)) e^{-ins} ds \right| \\ &< \varepsilon + (b - a)\varepsilon = (1 + b - a)\varepsilon. \end{aligned}$$

□

**Example 1.3.** We return to the functions discussed in Examples 1.1 and 1.2. Indeed,  $f_1(x) = x$  is smooth except at  $n\pi$ . According to Theorem 1.5, the series

$$2 \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} \sin nx$$

converges to  $x$  for every  $x \in (-\pi, \pi)$ . Taking  $x = \pi/2$ , we get the formula

$$\frac{\pi}{4} = 1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \dots.$$

On the other hand, we observed before that the series tend to 0 at  $x = \pm\pi$ . As  $f_1(\pi_+) = -\pi$  and  $f_1(\pi_-) = \pi$ , we have  $f_1(\pi_+) + f_1(\pi_-) = 0$ , which is in consistency with Theorem 1.5. In the second example,  $f_2(x) = x^2$  is continuous,  $2\pi$ -periodic. By Theorem 1.7, its Fourier series

$$\frac{\pi^2}{3} - 4 \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n^2} \cos nx$$

converges to  $x^2$  uniformly on  $[-\pi, \pi]$ . Taking  $x = 0$ , we obtain the formula

$$\frac{\pi^2}{12} = 1 - \frac{1}{4} + \frac{1}{9} - \frac{1}{16} + \dots.$$

So far we have been working on the Fourier series of  $2\pi$ -periodic functions. It is clear that the same results apply to the Fourier series of  $2T$ -periodic functions for arbitrary positive  $T$ .

We have shown the convergence of the Fourier series under some additional regularity assumptions on the function. But the basic question remains, that is, is the Fourier series of a continuous,  $2\pi$ -periodic function converges to itself? It turns out the answer

is negative. A not-so-explicit example can be found in Stein-Shakarchi and an explicit but complicated one was given by Fejér (see Zygmund “Trigonometric Series”). You may google for more. In fact, using the uniform boundedness principle in functional analysis, one can even show that “most” continuous functions have divergent Fourier series. The situation is very much like in the case of the real number system where transcendental numbers are uncountable while algebraic numbers are countable despite the fact that it is difficult to establish a concrete number is transcendental.

## Appendix I Series of Functions

This appendix serves to refresh your memory after the long hot summer.

A (real) sequence is a mapping  $\varphi$  from  $\mathbb{N}$  to  $\mathbb{R}$ . For  $\varphi(n) = a_n$ , we usually denote the sequence by  $\{a_n\}$  rather than  $\varphi$ . This is a convention. We say the sequence is convergent if there exists a real number  $a$  satisfying, for every  $\varepsilon > 0$ , there exists some  $n_0$  such that  $|a_n - a| < \varepsilon$  for all  $n, n \geq n_0$ . When this happens, we write  $a = \lim_{n \rightarrow \infty} a_n$ .

An (infinite) series is always associated with a sequence. Given a sequence  $\{x_n\}$ , set  $s_n = \sum_{k=1}^n x_k$  and form another sequence  $\{s_n\}$ . This sequence is the infinite series associated to  $\{x_n\}$  and is usually denoted by  $\sum_{k=1}^{\infty} x_k$ . The sequence  $\{s_n\}$  is also called the sequence of  $n$ -th partial sums of the infinite series. By definition, the infinite series is convergent if  $\{s_n\}$  is convergent. When this happens, we denote the limit of  $\{s_n\}$  by  $\sum_{k=1}^{\infty} x_k$ , in other words, we have

$$\lim_{n \rightarrow \infty} \sum_{k=1}^n x_k = \sum_{k=1}^{\infty} x_k.$$

So the notation  $\sum_{k=1}^{\infty} x_k$  has two meanings, first, it is the notation for an infinite series and, second, the limit of its partial sums (whenever it exists).

When the target  $\mathbb{R}$  is replaced by  $\mathbb{C}$ , we obtain a sequence or a series of complex numbers, and the above definitions apply to them after replacing the absolute value by the complex absolute value or modulus.

Let  $\{f_n\}$  be a sequence of real- or complex-valued functions defined on some non-empty  $E$  on  $\mathbb{R}$ . It is called convergent pointwisely to some function  $f$  defined on the same  $E$  if for every  $x \in E$ ,  $\{f_n(x)\}$  converges to  $f(x)$  as  $n \rightarrow \infty$ . Keep in mind that  $\{f_n(x)\}$  is sequence of real or complex numbers, so its convergence has a valid meaning. A more important concept is the uniform convergence. The sequence  $\{f_n\}$  is uniformly convergent to  $f$  if, for every  $\varepsilon > 0$ , there exists some  $n_0$  such that  $|f_n(x) - f(x)| < \varepsilon$  for all  $n \geq n_0$ . In notation  $f_n \rightrightarrows f$ . Equivalently, uniform convergence holds if, for every  $\varepsilon > 0$ , there exists some  $n_1$  such that  $\|f_n - f\|_{\infty} < \varepsilon$  for all  $n \geq n_1$ . Here  $\|f\|_{\infty}$  denotes the sup-norm of  $f$  on  $E$ .

An (infinite) series of functions is the infinite series given by  $\sum_{k=1}^{\infty} f_k(x)$  where  $f_k$  are defined on  $E$ . Its convergence and uniform convergence can be defined via its partial sums  $s_n(x) = \sum_{k=1}^n f_k(x)$  as in the case of sequences of numbers.

Among several criteria for uniform convergence, the following test is the most useful one.

**Weierstrass  $M$ -Test.** Let  $\{f_k\}$  be a sequence of functions defined on some  $E \subset \mathbb{R}$ . Suppose that there exists a sequence of non-negative numbers,  $\{a_k\}$ , such that

(a)  $|f_k(x)| \leq a_k$  for all  $k \geq 1$ , and

(b)  $\sum_{k=1}^{\infty} a_k$  is convergent.

Then  $\sum_{k=1}^{\infty} f_k$  converges uniformly and absolutely on  $E$ .

Also, the following “exchange theorem”.

**Exchange Theorem.** Let  $s_n = \sum_{k=1}^n f_k$  be uniformly convergent to  $\sum_{k=1}^{\infty} f_k$  on some  $E \subset \mathbb{R}$ . Then

(a)  $\sum_{k=1}^{\infty} f_k \in C(E)$  if  $f_k \in C(E)$  for all  $k$ .

(b) If  $E$  is an interval and  $f_k$ 's are differentiable with  $\sum_{k=1}^n f'_k \Rightarrow \sum_{k=1}^{\infty} f'_k$ , then  $\sum_{k=1}^{\infty} f_k$  is also differentiable and

$$\left( \sum_{k=1}^{\infty} f_k \right)' = \sum_{k=1}^{\infty} f'_k.$$

## Appendix II Sets of Measure Zero

Let  $E$  be a subset of  $\mathbb{R}$ . It is called of measure zero, or sometimes called a null set, if for every  $\varepsilon > 0$ , there exists a (finite or infinite) sequence of intervals  $\{I_k\}$  satisfying (1)  $E \subset \cup_{k=1}^{\infty} I_k$  and (2)  $\sum_{k=1}^{\infty} |I_k| < \varepsilon$ . (When the intervals are finite, the upper limit of the summation should be changed accordingly.) Here  $I_k$  could be an open, closed or any other interval and its length  $|I_k|$  is defined to be  $b - a$  where  $a \leq b$  are the endpoints of  $I_k$ .

The empty set is a set of measure zero from this definition. Every finite set is also null. For, let  $E = \{x_1, \dots, x_N\}$  be the set. For  $\varepsilon > 0$ , the intervals  $I_k = (x_1 - \varepsilon/(4N), x_k + \varepsilon/(4N))$  clearly satisfy (1) and (2) in the definition.

Next we claim that every countable set is also of measure zero. Let  $E = \{x_1, x_2, \dots\}$  be a countable set. We choose

$$I_k = \left( x_k - \frac{\varepsilon}{2^{k+2}}, x_k + \frac{\varepsilon}{2^{k+2}} \right) .$$

Clearly,  $E \subset \cup_{k=1}^{\infty} I_k$ . On the other hand,

$$\begin{aligned} \sum_{k=1}^{\infty} |I_k| &= \sum_{k=1}^{\infty} \frac{\varepsilon}{2^{k+1}} \\ &= \frac{\varepsilon}{2} \\ &< \varepsilon . \end{aligned}$$

We conclude that every countable set is a null set.

There are uncountable sets of measure zero. For instance, the Cantor set which plays an important role in analysis, is of measure zero. Here we will not go into this.

The same trick in the above proof can be applied to the following situation.

**Proposition A.1.** *The union of countably many null sets is a null set.*

*Proof.* Let  $E_k, k \geq 1$ , be sets of measure zero. For  $\varepsilon > 0$ , there are intervals satisfying  $\{I_j^k\}, E_k \subset \cup_j I_j^k$ , and  $\sum_j |I_j^k| < \varepsilon/2^k$ . It follows that  $E \equiv \cup_k E_k \subset \cup_{j,k} I_j^k = \cup_k \cup_j I_j^k$  and

$$\sum_k \sum_j |I_j^k| < \sum_k \frac{\varepsilon}{2^k} = \varepsilon .$$

□

The concept of a null set comes up naturally in the theory of Riemann integration. A theorem of Lebesgue asserts that a bounded function is Riemann integrable if and only if its discontinuity set is null. The following result is used in the uniqueness assertion on Fourier series. I provide a proof here, but you may just take it for granted.

**Proposition A.2.** *Let  $f$  be a non-negative integrable function on  $[a, b]$ . Then  $\int_a^b f = 0$  if and only if  $f$  is equal to 0 except on a null set. Consequently, two integrable functions  $f, g$  satisfying*

$$\int_a^b |f - g| = 0,$$

*if and only if  $f$  is equal to  $g$  except on a null set.*

*Proof.* Let  $f$  be a non-negative integrable function satisfying  $\int_a^b f = 0$ . We set, for each  $k \geq 1$ ,  $A_k = \{x \in [a, b] : f(x) > 1/k\}$ . It is clear that

$$\{x : f(x) > 0\} = \bigcup_{k=1}^{\infty} A_k .$$

By Proposition A.1., it suffices to show that each  $A_k$  is null. Thus let us consider  $A_{k_0}$  for a fixed  $k_0$ . Recall from the definition of Riemann integral, for every  $\varepsilon > 0$ , there exists a partition  $a = x_1 < x_2 < \cdots < x_n = b$  such that

$$0 \leq \sum_{k=1}^{n-1} f(z_k)|I_k| = \left| \sum_{k=1}^{n-1} f(z_k)|I_k| - \int_a^b f \right| < \frac{\varepsilon}{k_0},$$

where  $I_k = [x_k, x_{k+1}]$  and  $z_k$  is an arbitrary tag in  $[x_k, x_{k+1}]$ . Let  $\{k_1, \dots, k_m\}$  be the index set for which  $I_{k_j}$  contains a point  $z_{k_j}$  from  $A_{k_0}$ . Choosing the tag point to be  $z_{k_j}$ , we have  $f(z_{k_j}) > 1/k_0$ . Therefore,

$$\frac{1}{k_0} \sum_{k_j} |I_{k_j}| = \sum_{k_j} f(z_{k_j})|I_{k_j}| \leq \sum_{k=1}^{n-1} f(z_k)|I_k| < \frac{\varepsilon}{k_0},$$

so

$$\sum_{k_j} |I_{k_j}| < \varepsilon.$$

We have shown that  $A_{k_0}$  is of measure zero.

Conversely, let  $I_k, k = 1, \dots, n$ , be a partition of  $[a, b]$ . Since the non-zero set of  $f$  is of measure zero, from each subinterval  $I_k$  one can pick a tag point  $z_k$  such that  $f(z_k) = 0$ . As a result, the Riemann sum  $\sum_{k=1}^n f(z_k)|I_k| = 0$ . As the Riemann sums approach the Riemann integrable as the length of partitions tend to 0, we conclude that  $\int_a^b f = 0$ .

□

A property holds **almost everywhere** if it holds except on a null set. For instance, this proposition asserts that the integral of a non-negative function is equal to zero if and only if it vanishes almost everywhere.

**Comments on Chapter 1.** According to the development of Undergraduate Analysis, Fourier series should be the topic right after power series. However, it was not presented in MATH2060 due to lack of time.

Historically, the relation (2.2) comes from a study on the one-dimensional wave equation

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

where  $u(x, t)$  denote the displacement of a string at the position-time  $(x, t)$ . Around 1750, D'Alembert and Euler found that a general solution of this equation is given by

$$f(x - ct) + g(x + ct)$$

where  $f$  and  $g$  are two arbitrary twice differentiable functions. However, D. Bernoulli found that the solution could be represented by a trigonometric series. These two different ways of representing the solutions led to a dispute among the mathematicians at that time, and it was not settled until Fourier gave many convincing examples of representing functions by trigonometric series in 1822. His motivation came from heat conduction. After that, trigonometric series have been studied extensively and people call it Fourier series in honor of the contribution of Fourier. Nowadays, the study of Fourier series has matured into a branch of mathematics called harmonic analysis. It has equal importance in theoretical and applied mathematics, as well as other branches of natural sciences and engineering.

In some books the Fourier series of a function is written in the form

$$\frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx),$$

instead of

$$a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx),$$

so that the formula for  $a_0$  is the same as the other  $a_n$ 's (see (2.1)). However, our notation has the advantage that  $a_0$  has a simple meaning, i.e., it is the average of the function over a period.

Concerning the convergence of a Fourier series to its function, we point out that an example of a continuous function whose Fourier series diverges at some point can be found in Stein-Sharachi. More examples are available by googling. The classical book by A. Zygmund, "Trigonometric Series" (1959) reprinted in 1993, contains most results before 1960. After 1960, one could not miss to mention Carleson's sensational work in 1966, whose result implies that the Fourier series of every function in  $R_{2\pi}$  converges to the function itself almost everywhere.

The aim of this chapter is to give an introduction to Fourier series. It will serve the purpose if your interest is aroused and now you consider to take our course on Fourier analysis in the future.